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COMPRESSOR FREE HYBRID (SOLID/GAS) FS SYSTEM

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Photochemically Driven Active Media: New Strategy in the Development of Ultra-high Power Fs Systems

OUTLINE

- **Main properties of broadband photochemically driven active media**
- **Architecture of a hybride multiterawatt systems**
- **XeF(C-A)-amplifier optically driven by the 172 nm radiation from an e-beam pumped flash-lamp**
- **Recompression of fs pulses**
- **Conclusions**

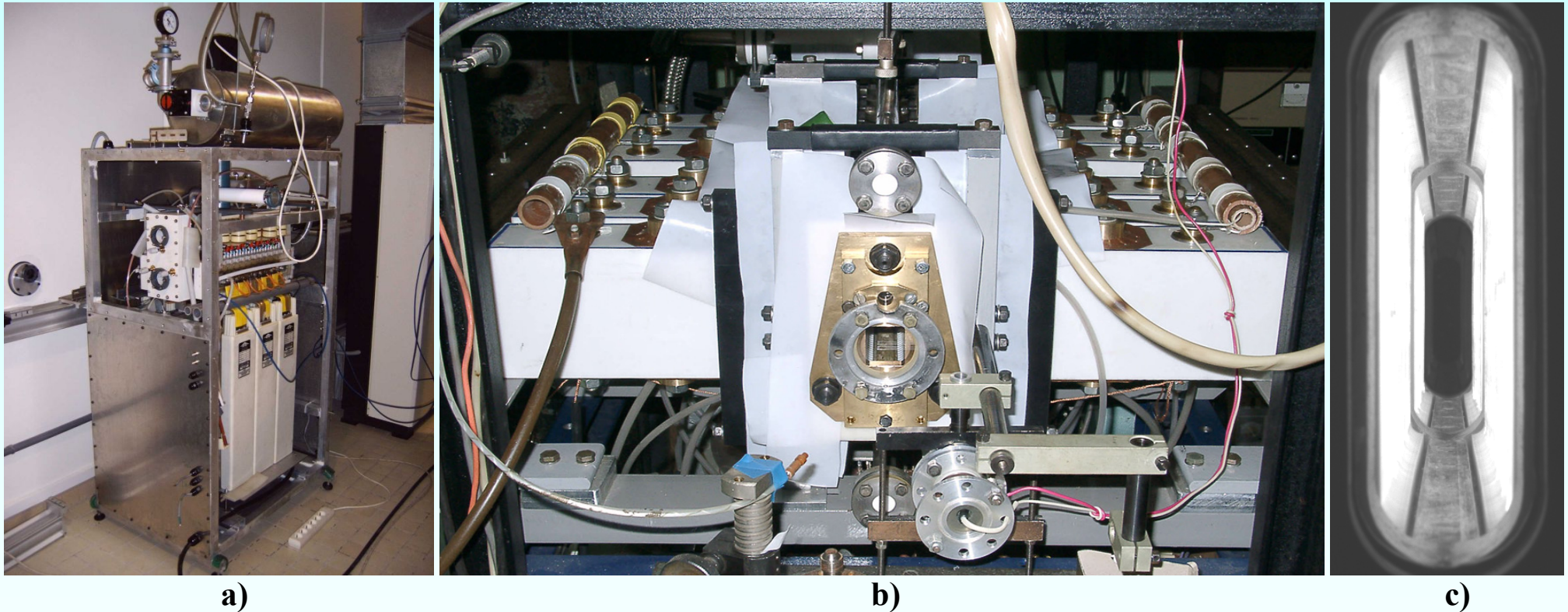
Advantages of gaseous optically pumped active media: Low optical nonlinearity → CPA in subps time domain
 Scalability to very large volumes
 Low cost of realization
 High temporal contrast ($>10^{10}$)

***Photochemically driven broadband active media
 for the fs optical pulse amplification***

Transition	XeF(C-A) /1/	Kr ₂ F (4 ² Γ-1,2 ² Γ) /2/	Xe ₂ Cl (4 ² Γ-1,2 ² Γ) /3/
λ_{\max} , nm	474	405	485
$\Delta\lambda$, nm	60-100	80	100
τ_{lim} , fs	8-12	7	8
τ_{sp} , ns	100	181	245
σ_{st} , cm ²	10^{-17}	2.3×10^{-18}	2.8×10^{-18}
ϵ_{sat} , J/cm ²	0.05	0.2	0.15
I, TW/cm ² ($\tau = 25$ fs)	2	8	6

1. L.D.Mikheev, D.B.Stavrovskii, V.S.Zuev: J. Rus. Las. Res., v.16, 427 (1995)
2. N.G.Basov, V.S.Zuev, A.V.Kanaev, L.D.Mikheev, D.B.Stavrovskii: Kvantovaya Elektron.(Moscow) v.7, 2660 (1980)
3. N.G.Basov, V.S.Zuev, A.V.Kanaev, L.D.Mikheev: Kvantovaya Elektron. (Moscow) v.12, 1954 (1985)

XeF(C-A) amplifier pumped by radiation from a surface discharge



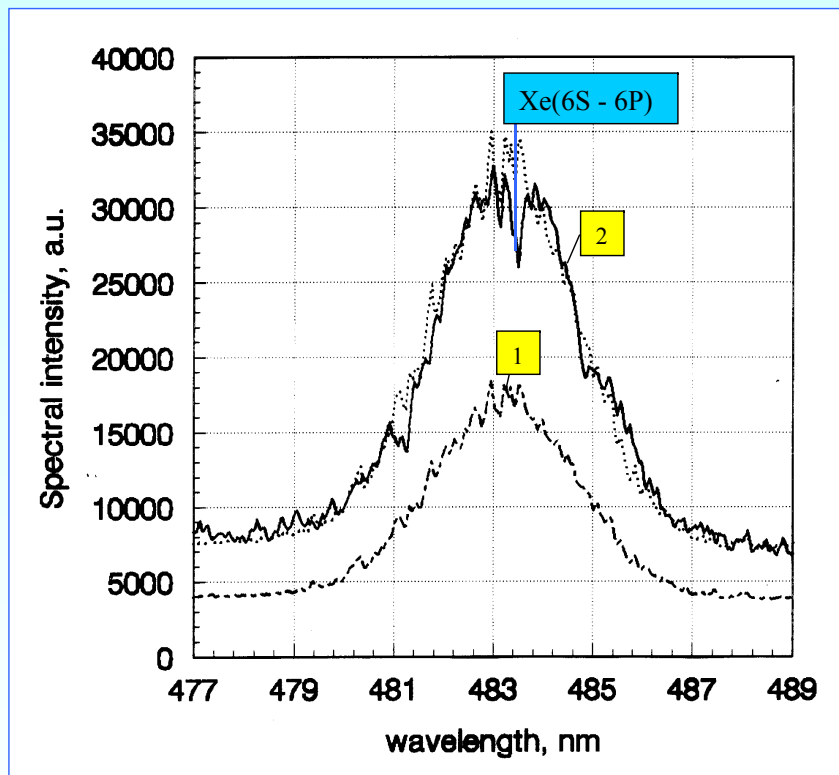
Photos of the photolytical XeF(C-A) amplifiers: a) LP3 of Marseille Univ.(active volume: $5 \times 18 \times 40 \text{ cm}^3$); b) P.N.Lebedev Inst. (active volume: $3 \times 11 \times 50 \text{ cm}^3$); c) XeF(C-A) amplifier viewed from its front when surface discharge is initiated.

Small-signal gain: $2 \times 10^{-3} \text{ cm}^{-1}$
Total amplification factor: 10^2

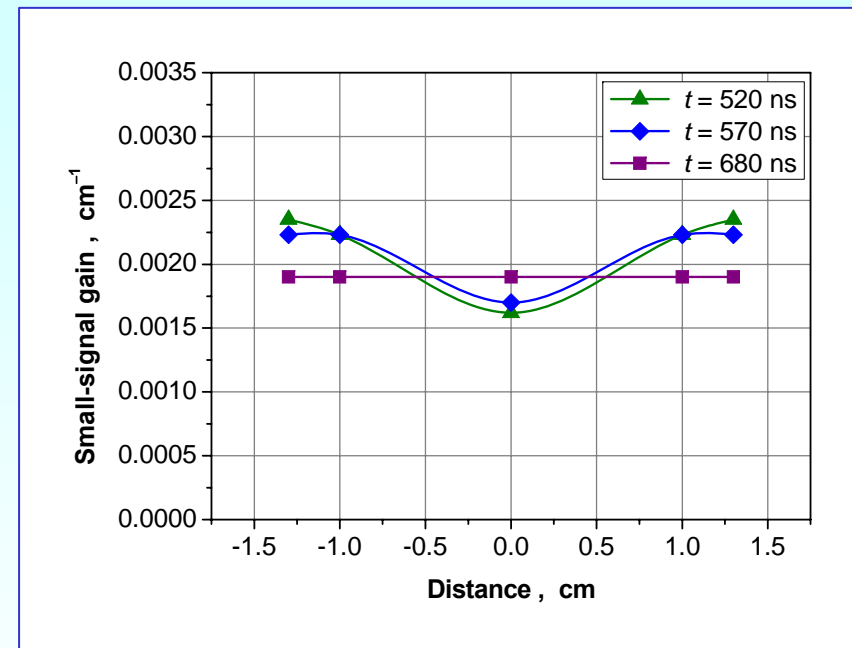
Surface discharge is promising up to 100 TW

XeF(C-A) amplifier photochemically driven by radiation from a surface discharge. Experimental results.

Total small signal gain of the multipass amplifier 10^2

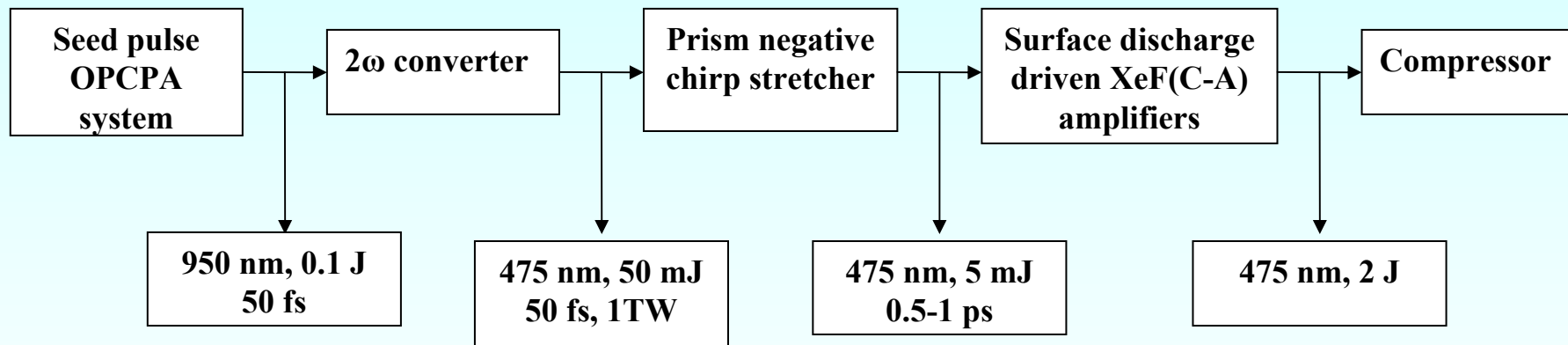


Spectra of fs pulse before (1) and after (2) amplification



Gain distributions versus the distance from the central plane of the amplifier at different instants of time

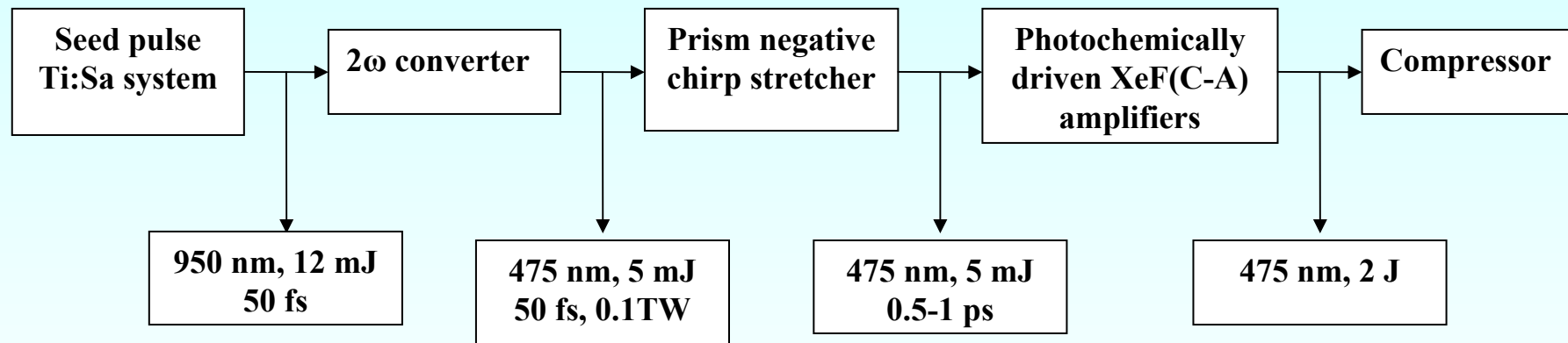
Architecture of a 10 Tw hybrid fs XeF(C-A) system at LP3 (Marseille University)



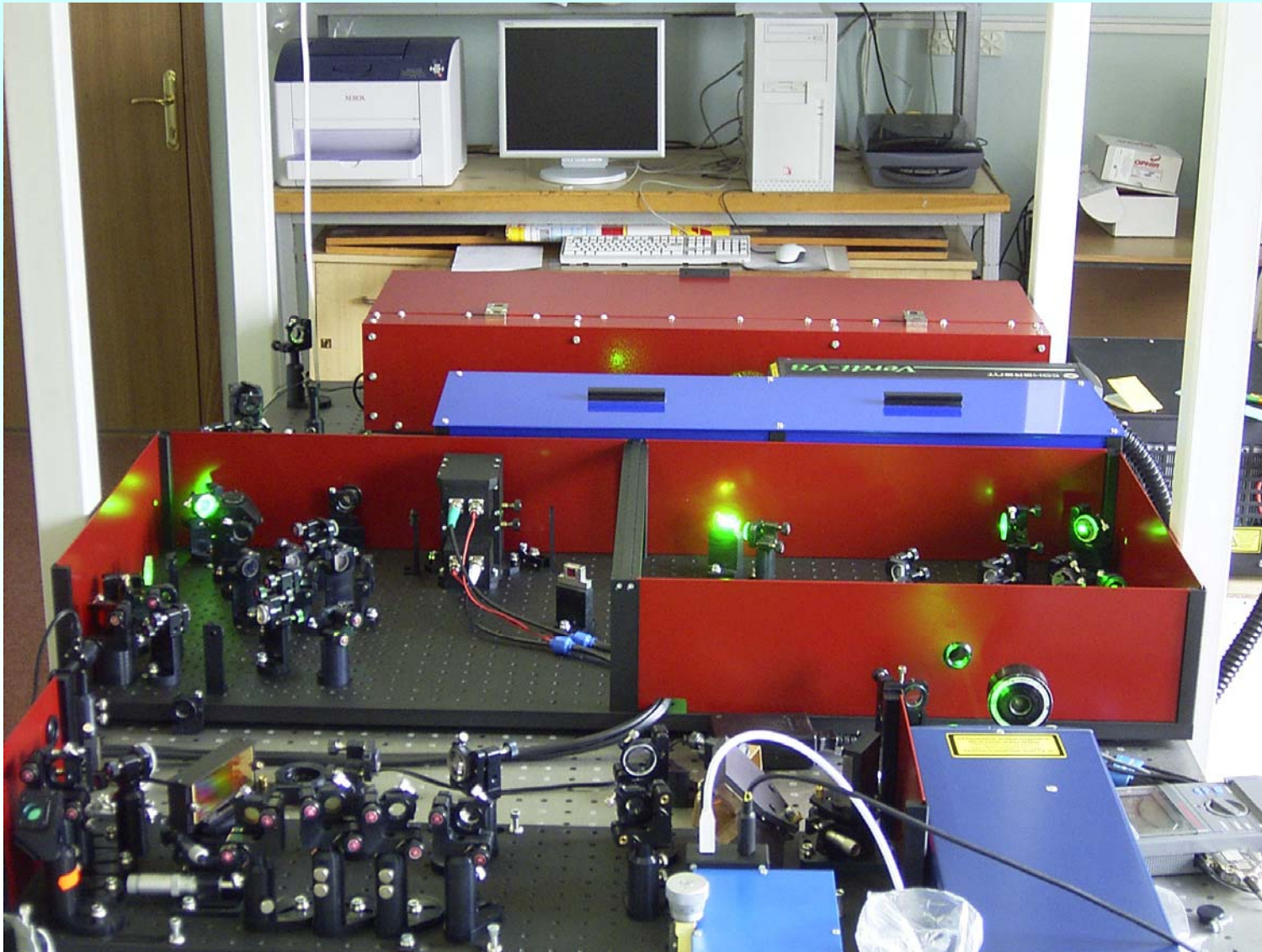
R.Clady, G.Coustillier, M.Gastaud, M.Sentis, P.Spiga,
V.Tcheremiskine, O.Uteza, L.D.Mikheev, V.Mislavskii,
J.P.Chambaret, G.Chériaux.

Appl. Phys. B, **82**, 347-358 (2006).

Architecture of a 100 TW hybrid fs XeF(C-A) system at P.N.Lebedev Inst. (Moscow)



Ti:Sa front end (Avesta Project Ltd)



$$\lambda_{\omega} = 950 \text{ nm}$$

$$\lambda_{2\omega} = 475 \text{ nm}$$

$$\tau = 45\text{-}50 \text{ fs}$$

$$E = 5 \text{ mJ}$$

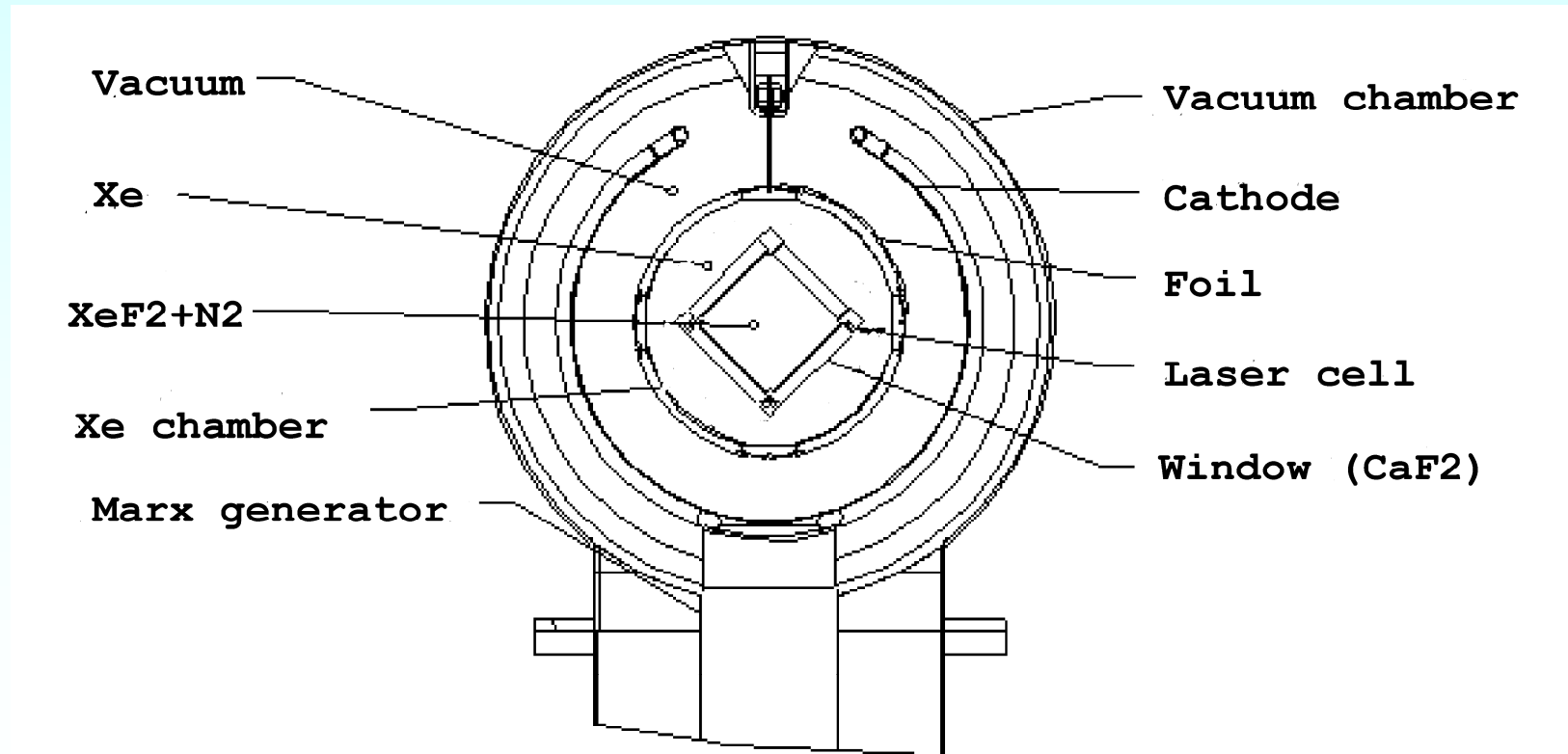
$$f = 10 \text{ Hz}$$

Final XeF(C-A) amplifier photochemically driven by 172 nm radiation from an e-beam pumped Xe converter

Active medium length - 1.2 m, clear aperture - 12 cm.

E-beam: $I=80 \text{ } \mu\text{A}$, $U_e= 420 \text{ keV}$, pulse-width – 400 ns.

Xe_2 fluorescence efficiency related to e-beam energy is 30-40%

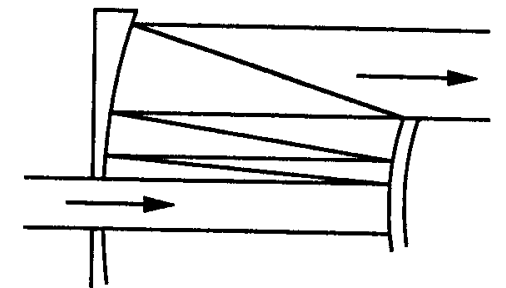


XeF(C-A) amplifier (IHCE, Tomsk)



Experiment:
Small signal gain
 $2.6 \times 10^{-3} \text{ cm}^{-1}$

Theory:
Small signal gain
 $5 \times 10^{-3} \text{ cm}^{-1}$
 $E_{\text{stor}} = 6 \text{ J}$
 $E_{\text{out}} = 2\text{-}2.5 \text{ J}$
 $P_{\text{out}} = 50 \text{ TW}$
 $\tau = 50 \text{ fs}$



Recompression of downchirped fs pulses in bulk materials (physical insight)

PW systems – recompression in bulk materials in linear regime

Is it possible to do it in output window (nonlinear interaction)?

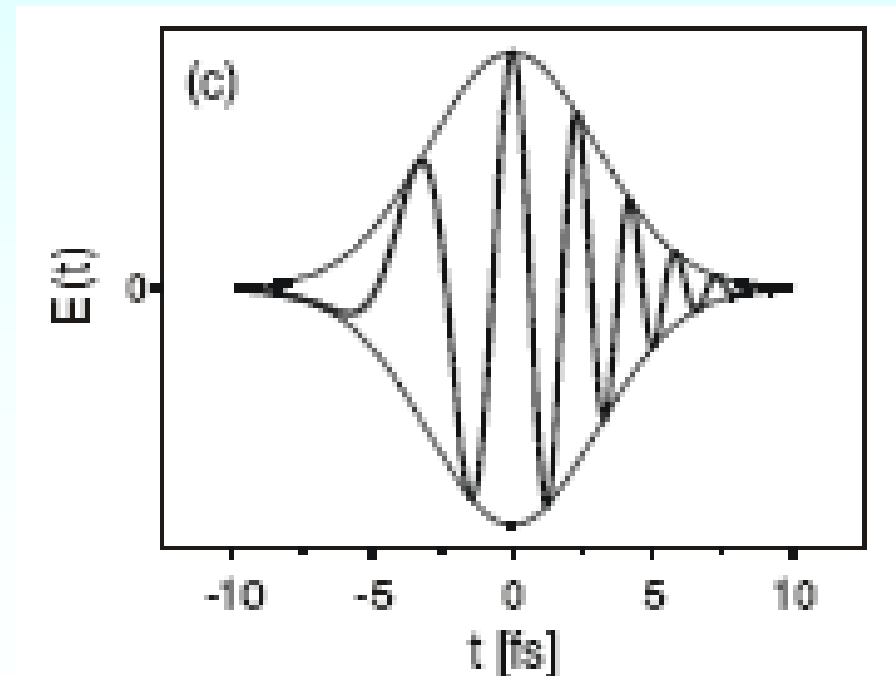
Gaseous active media $\Rightarrow W_{\text{out}} \sim (1-4) \text{ TW/cm}^2 \ll W_{\text{dam}}$ (bulk glass, gas)

Self-phase modulation

$$\Delta\omega \sim (-dn/dt)$$

Kerr effect: $dn/dt > 0, \Delta\omega < 0$

Electron plasma: $dn/dt < 0, \Delta\omega > 0$



Recompression of downchirped fs pulses in bulk materials (physical insight)

$$\Delta n = n_2 I h\nu + \sqrt{n_0^2 - \frac{\omega_p^2}{\omega^2}} - n_0 \approx n_2 I h\nu - \frac{1}{2n_0} \frac{\omega_p^2}{\omega^2} = n_2 I h\nu - \frac{1}{2n_0} \frac{\rho^e}{\rho_{\text{crit}}^e}$$

$n_2 = 3.5 \times 10^{-16} \text{ cm}^2/\text{W}$ is the fused silica nonlinear refractive index,
 I [phot/cm²s] is the intensity of the electric field,

$\omega_p = \sqrt{\frac{\rho^e e^2}{\epsilon_0 n_0^2 m}}$ is the plasma frequency, ω is the carrier frequency,

$m = 0.635m_e$ denotes the reduced mass of the electron and the hole,

ρ^e stands for the electron plasma density,

$\rho_{\text{crit}}^e = \epsilon_0 n_0^2 m \omega^2 / e^2 \approx 4.2 \times 10^{21} \text{ cm}^{-3}$ is the critical plasma density at 480 nm

Recompression of downchirped fs pulses in bulk materials (physical insight)

Electron plasma generation in fused silica

($U_i = 7.5-9$ eV, $h\nu$ (480 nm) = 2.5 eV, $U_i / h\nu=3-4$, $W_{out} \sim (1-4)$ TW/cm²):

1. Multiphoton ionization (MPI)
2. Tunneling ionization – minor due to $\gamma=\omega(mU_i)^{1/2}/eE \gg 1$
3. Avalanche ionization – minor?

MPI: $\rho^e \approx \sigma_k I^k N_a \tau_p$, $k=3,4$

$\sigma_3 = 6 \times 10^{-81}$ cm⁶c² (S.C.Jones et al., Opt.Eng. 28, 1039 (1989))

$\sigma_4 = 2 \times 10^{-114}$ cm⁸c³ (S.C.Jones et al., Opt.Eng. 28, 1039 (1989))

$N_a = 2.2 \times 10^{22}$ cm⁻³, $\tau_p = 30$ fs \ll electron recombination time of 150 fs.

$\Delta n = 0$ at $W = 4.5$ TW/cm² ($k=4$) or $W = 0.7$ TW/cm² ($k=3$)

$\rho^e \approx 2 \times 10^{19}$ cm³ ($k=4$) or 3.2×10^{18} cm³ ($k=3$)

Recompression of downchirped fs pulses in bulk fused silica (numerical simulation)

Downchirped Gaussian pulse:

$$A|_{z=0} = A_0 \exp \left[-(\ln 2) \left(\left(\frac{t-t_0}{\tau} \right)^2 + \frac{r^2}{2a^2} \right) - \frac{i\chi(t-t_0)^2}{2} \right]$$

traveling through a piece of fused silica with normal dispersion. Combined nonlinear (generalized) Schrödinger equation in dimensionless variables:

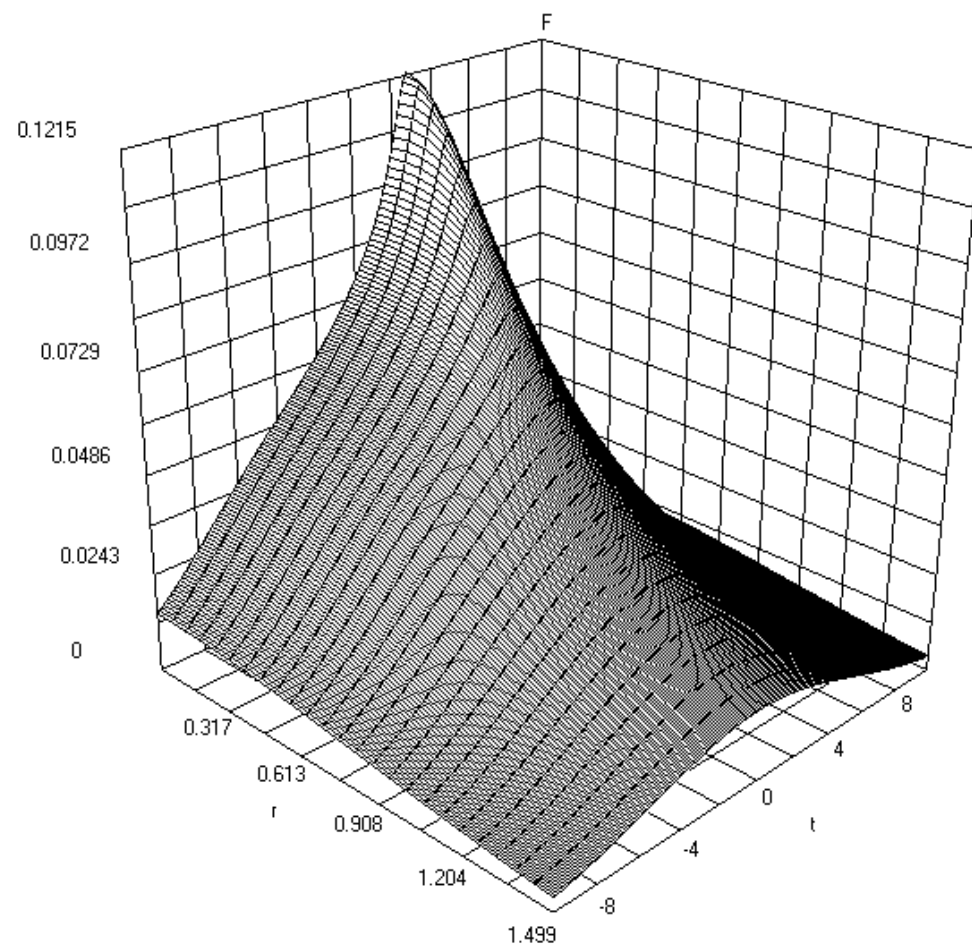
$$\frac{\partial A}{\partial z} + iD_r \Delta_r A + iD \frac{\partial^2 A}{\partial t^2} + i(\alpha |A|^2 - \beta\rho)A + \gamma \frac{\partial}{\partial t} ((\alpha |A|^2 - \beta\rho)A) = 0$$

$$0 < z \leq L_z, \quad 0 < t < L_t, \quad 0 < r < R$$

$$\frac{\partial \rho}{\partial t} + \frac{\rho}{\tau_{rel}} = \delta(1 - \rho) |A|^{2n}, \quad n = 3$$

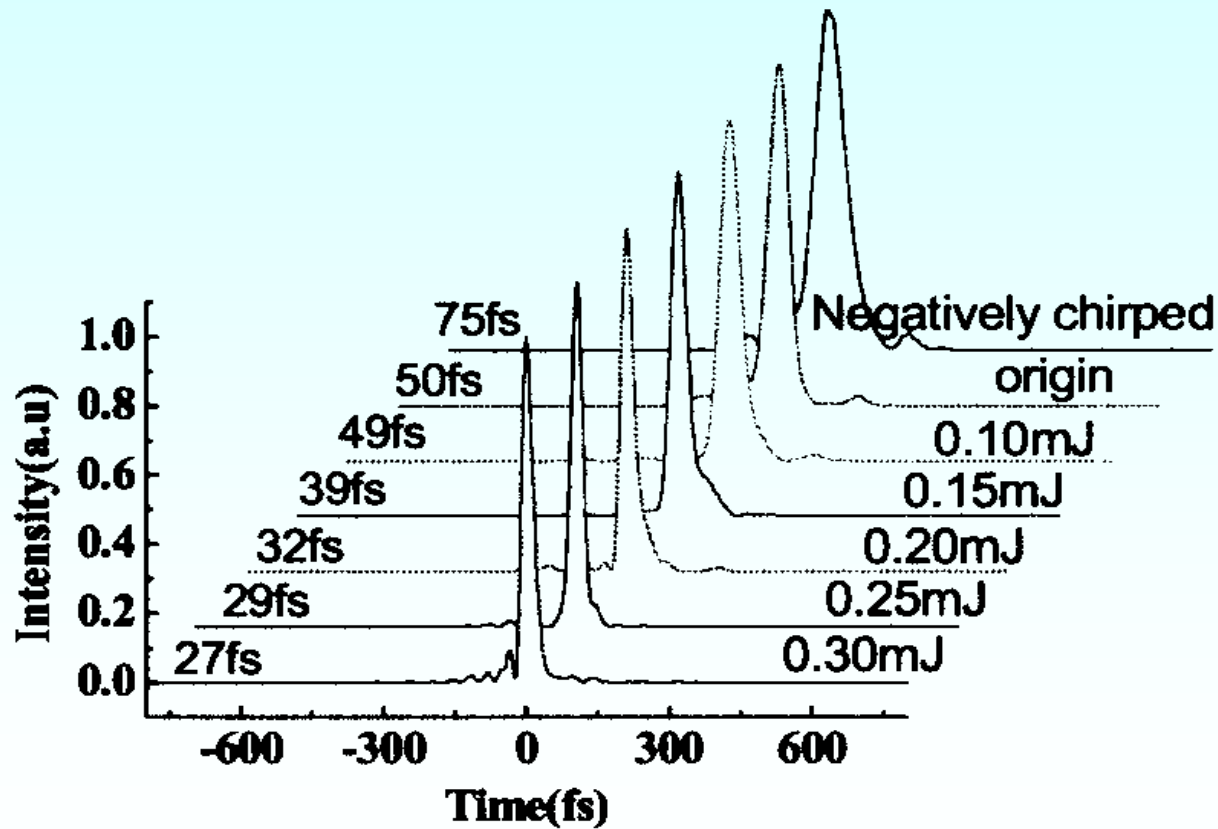
Where

$$D_r = L^{-1}_{diffr}, \quad D = L^{-1}_{disp}, \quad \alpha = L^{-1}_{sf}, \quad \gamma = 2(\omega_0 \tau_{pulse})^{-1}, \quad \delta = \sigma^{(3)} N I^3, \quad \beta = k_0 N_a / 2\rho_{cr}, \quad W = 1.6 \text{ TW/cm}^2$$



Nonlinear compression of negatively chirped fs pulses in BK7 (experiment)

Jun Liu et al. Opt.Express, v.14, 979 (2006)



BK-7

$\lambda = 800 \text{ nm}$

1.7 TW/cm^2

Conclusions

- Optimistic experimental and theoretical results obtained show that realization of compressor-free hybrid fs system is expected to be feasible.
- Preparation work is now in progress to demonstrate 50 TW peak power in the hybrid fs system which is being built at P.N.Lebedev Institute.